

Reciprocal Lattice

Recall: plane waves of form $e^{i\vec{k}\cdot\vec{r}}$, for arbitrary \vec{k} & \vec{r} .

Now restrict \vec{k} st. the plane wave reflects the symmetry of the direct lattice $\Rightarrow \vec{K}$

That is:

$$e^{i\vec{K}\cdot\vec{r}} = e^{i\vec{K}\cdot(\vec{r}+\vec{R})}$$

Since \vec{R} represent any pt. of a Bravais Lattice

$$\vec{R} = n_1\vec{a}_1 + n_2\vec{a}_2 + n_3\vec{a}_3$$

$$\therefore \boxed{e^{i\vec{K}\cdot\vec{R}} = 1 \text{ for all } \vec{R}}$$

This set of \vec{K} define the pts of the reciprocal lattice of the Bravais lattice defined by the set of \vec{R} .

Reciprocal lattice: The set of all \vec{K} in k -space (momentum space) that give plane waves which have the symmetry of a Bravais lattice, is known as the reciprocal lattice.

Note: "Direct lattice" refers to a real-space Bravais lattice

- the dft used above is equivalent to the dft of a Fourier pair:

i.e. the direct lattice & reciprocal lattice are Fourier transforms of each other!

Describing the R.L.:

$$e^{i\vec{k} \cdot \vec{R}} = 1$$

$$\therefore \vec{k} \cdot \vec{R} = 2\pi m, \quad m \rightarrow \text{integer}$$

$$\Leftrightarrow \vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$$

$$\text{write } \vec{k} = m_1 \vec{b}_1 + m_2 \vec{b}_2 + m_3 \vec{b}_3 \quad (m_i \text{ arbitrary for now})$$

and define:

$$\vec{b}_1 = \frac{2\pi}{\Omega_0} (\vec{a}_2 \times \vec{a}_3)$$

$$\Omega_0 = \vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)$$

Volume of direct lattice primitive cell.

$$\vec{b}_2 = \frac{2\pi}{\Omega_0} (\vec{a}_3 \times \vec{a}_1)$$

$$\vec{b}_3 = \frac{2\pi}{\Omega_0} (\vec{a}_1 \times \vec{a}_2)$$

$$\text{Note: } \vec{b}_i \cdot \vec{a}_j = 2\pi \delta_{ij}$$

$$\delta_{ij} = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$$

$$\therefore \bar{K} \cdot \bar{R} = 2\pi(n_1 m_1 + n_2 m_2 + n_3 m_3)$$

$\therefore n_1 m_1 + n_2 m_2 + n_3 m_3 = \text{integer}$, for any choice of n_i

$\therefore m_i$ also integers

\therefore The reciprocal lattice is defined by the pts:

$$\bar{K} = m_1 \bar{b}_1 + m_2 \bar{b}_2 + m_3 \bar{b}_3$$

where m_i are integers & \bar{b}_i are vectors (defined before) all not in the same plane.....

This is the defn of a Bravais Lattice

The reciprocal lattice of a B.L. is also a B.L. (But not the same B.L.).

You can repeat this process to calc. the reciprocal of the reciprocal lattice. This yields the original direct lattice.

\hookrightarrow Not surprising if you consider \bar{K} & \bar{R} as Fourier transforms.

Direct lattice

$$\bar{a}_1, \bar{a}_2, \bar{a}_3$$

Reciprocal lattice

$$\bar{b}_1, \bar{b}_2, \bar{b}_3$$

Both Bravais Lattices

$$\Omega_0 = \bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3)$$

$$\tilde{\Omega}_0 = \bar{b}_1 \cdot (\bar{b}_2 \times \bar{b}_3)$$

$$\tilde{\Omega}_0 = \bar{b}_1 \cdot (\bar{b}_2 \times \bar{b}_3)$$

$$= \frac{2\pi}{\Omega_0} (\bar{a}_2 \times \bar{a}_3) \cdot (\bar{b}_2 \times \bar{b}_3)$$

$$= \frac{2\pi}{\Omega_0} \bar{a}_2 \cdot (\bar{a}_3 \times \bar{b}_2 \times \bar{b}_3)$$

$$= \frac{2\pi}{\Omega_0} \bar{a}_2 \cdot \left[\underbrace{\bar{b}_2 (\bar{a}_3 \cdot \bar{b}_3)}_{2\pi \bar{b}_2} - \underbrace{\bar{b}_3 (\bar{a}_3 \cdot \bar{b}_2)}_0 \right]$$

$$= \frac{2\pi}{\Omega_0} 2\pi (\bar{a}_2 \cdot \bar{b}_2)$$

$$= \frac{8\pi^3}{\Omega_0}$$

Volume of direct primitive cell Ω_0

Volume of reciprocal prim. cell $\tilde{\Omega}_0$

$$\tilde{\Omega}_0 = \frac{8\pi^3}{\Omega_0}$$

Examples:

Simple cubic:

$$\bar{a}_1 = a \hat{x} \quad \bar{a}_2 = a \hat{y} \quad \bar{a}_3 = a \hat{z}$$

calculate \bar{b}_i :

$$\bar{b}_3 = \Omega_0 = \bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3) = a^3$$

$$\bar{b}_1 = \frac{2\pi}{\Omega_0} (\bar{a}_2 \times \bar{a}_3) = \frac{2\pi}{a^3} a^2 \hat{x}$$

$$\bar{b}_1 = \frac{2\pi}{a} \hat{x}$$

$$\bar{b}_2 = \frac{2\pi}{a} \hat{y}$$

$$\bar{b}_3 = \frac{2\pi}{a} \hat{z}$$

These are the prim.
vectors for a
sc Bravais Lattice
w sides $\frac{2\pi}{a}$

$$\text{FCC: } \bar{a}_1 = \frac{a}{2} (\hat{y} + \hat{z}), \quad \bar{a}_2 = \frac{a}{2} (\hat{z} + \hat{x}), \quad \bar{a}_3 = \frac{a}{2} (\hat{x} + \hat{y})$$

$$\Omega_0 = \bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3) = \frac{a^3}{4}$$

$$\bar{b}_1 = \frac{2\pi}{\Omega_0} (\bar{a}_2 \times \bar{a}_3) = \frac{8\pi}{a^3} \left(\frac{a^2}{4}\right) (-\hat{x} + \hat{y} + \hat{z})$$

$$= \frac{2\pi}{a} (\hat{y} + \hat{z} - \hat{x})$$

$$\bar{b}_2 = \frac{8\pi}{a^3} (\bar{a}_3 \times \bar{a}_1) = \frac{2\pi}{a} (\hat{z} + \hat{x} - \hat{y})$$

$$\bar{b}_3 = \frac{8\pi}{a^3} (\bar{a}_1 \times \bar{a}_2) = \frac{2\pi}{a} (\hat{x} + \hat{y} - \hat{z})$$

rewrite: $\frac{2\pi}{a} = \frac{4\pi}{a} \cdot \frac{1}{2}$

$$\therefore \left. \begin{aligned} \bar{b}_1 &= \frac{4\pi}{a} \cdot \frac{1}{2} (\hat{y} + \hat{z} - \hat{x}) \\ \bar{b}_2 &= \frac{4\pi}{a} \cdot \frac{1}{2} (\hat{z} + \hat{x} - \hat{y}) \\ \bar{b}_3 &= \frac{4\pi}{a} \cdot \frac{1}{2} (\hat{x} + \hat{y} - \hat{z}) \end{aligned} \right\} \begin{array}{l} \text{prim. vectors} \\ \text{for BCC} \\ \text{with cube} \\ \text{side length} \\ \frac{4\pi}{a} \end{array}$$

check: FCC: $\Omega_0 = \frac{a^3}{4} \rightarrow \text{side} = a.$

$$\text{From: } \hat{\Omega}_0 = \frac{8\pi^3}{\Omega_0} = \frac{32\pi^3}{a^3}$$

BCC: side = $\frac{4\pi}{a}$

$$\Omega = \left(\frac{4\pi}{a}\right)^3 \times \frac{1}{2} \quad (2 \text{ pts per unit cube})$$

$$= \frac{64\pi^3}{a^3} \cdot \frac{1}{2} = \frac{32\pi^3}{a^3} \quad \checkmark$$

So recip. of FCC is BCC. Of course \therefore recip of BCC is FCC!!

Primitive unit cell in reciprocal space:

→ since reciprocal lattice is also a Bravais lattice the primitive cells are similar to what we have already discussed

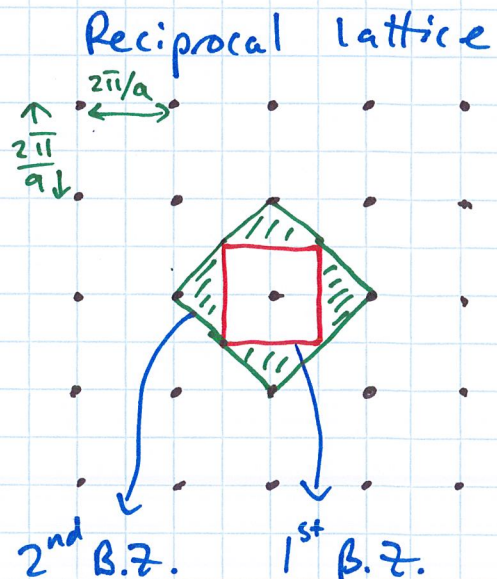
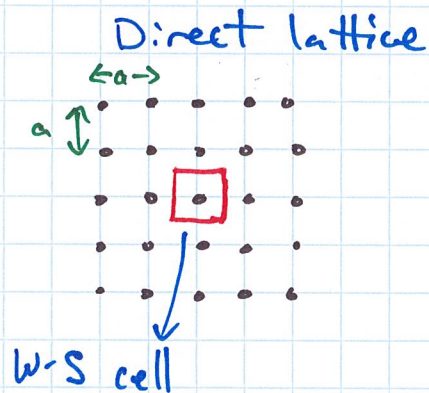
→ terminology is different though!

Direct lattice → Wigner-Seitz cell

Reciprocal lattice → First Brillouin Zone

→ geometrically identical

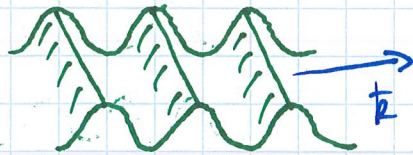
eg. square lattice



→ each B.Z. has the same total area/volume.

Lattice planes & Miller indices

plane wave: $e^{i\vec{k}\cdot\vec{r}}$



- travelling in direction \vec{k}

- wave is constant in plane \perp to \vec{k} .

- Each \vec{k} in the reciprocal lattice represents a plane (set of planes) in the direct lattice

→ the reverse is also true

Take arbitrary $\vec{k} = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$

real-space vector \vec{r} lies in a specific plane

if:

$$e^{i\vec{k}\cdot\vec{r}} = \text{constant}$$

or $\vec{k}\cdot\vec{r} = A$ (A is a const.)
for a given plane

where does this plane cross the real space axes \vec{a}_i ?

Define these pts: $x_1\vec{a}_1, x_2\vec{a}_2, x_3\vec{a}_3 = \vec{r}$

Then for each of these \vec{r} :

$$\vec{k}\cdot\vec{r} = A$$

$$\text{ie: } \bar{K} \cdot (x_i \bar{a}_i) = A$$

$$\text{Since } \bar{K} = h \bar{b}_1 + k \bar{b}_2 + l \bar{b}_3$$

$$\therefore \bar{K} \cdot \bar{a}_1 = 2\pi h \quad \text{etc. . .}$$

$$\therefore \bar{K} \cdot (x_1 \bar{a}_1) = 2\pi h x_1 = A$$

$$\text{or } x_1 = \frac{A}{2\pi h}$$

$$\text{likewise: } x_2 = \frac{A}{2\pi k} \quad x_3 = \frac{A}{2\pi l}$$

The real-space intercepts of the plane defined by h, k, l are inversely related to h, k, l .

h, k, l define a Miller plane called Miller indices.

h, k, l are defined such that:

$$h : k : l = \frac{1}{x_1} : \frac{1}{x_2} : \frac{1}{x_3}$$

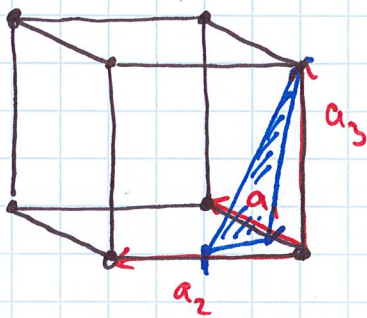
and h, k, l have no common factors
(remember they are also integers)

e.g. on $(hkl) = (421)$ plane intercepts

the \bar{a}_1 axis @ $\frac{1}{h} = \frac{1}{4} \bar{a}_1$

a_2 axis @ $\frac{1}{k} = \frac{1}{2} \bar{a}_2$

a_3 axis @ $\frac{1}{l} = \frac{1}{1} \bar{a}_3$



→ Since any integer multiple of \bar{K} is also a R.L. pt and is \parallel to \bar{K} , (hkl) actually represent an infinite set of \parallel planes

→ the periodicity of these planes is $\frac{2\pi}{|\bar{K}|}$

i.e. planes in the family (hkl)

where $\bar{K} = h\bar{b}_1 + k\bar{b}_2 + l\bar{b}_3$

are separated by a distance of $d = \frac{2\pi}{|\bar{K}|}$

Important notes about Miller indices

- Miller plane: $(h \ k \ l)$

$$\bar{K} = h \bar{b}_1 + k \bar{b}_2 + l \bar{b}_3 \quad \text{is R.L. vector}$$

- The real-space plane is \perp to \bar{K}

- h, k, l are integers w/ no common factors

ie. $(4 \ 6 \ 2) \Rightarrow (2 \ 3 \ 1) \checkmark$

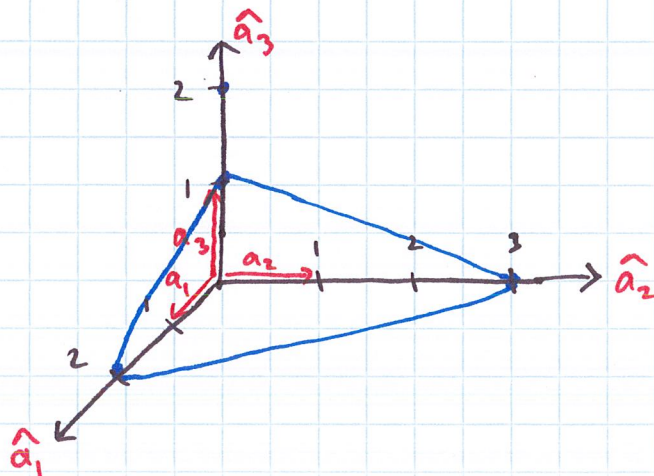
- real space distance between family of planes is:

$$d = \frac{2\pi}{|\bar{K}|}$$

- the plane intercepts the realspace axes $a_{1,2,3}$
@ pts $x_{1,2,3}$

where $\frac{1}{h} : \frac{1}{k} : \frac{1}{l} = \frac{1}{x_1} : \frac{1}{x_2} : \frac{1}{x_3}$

e.g. find the Miller indices for the plane
shown below:



Intercepts: $x_1 = 2$
 $x_2 = 3$
 $x_3 = 1$

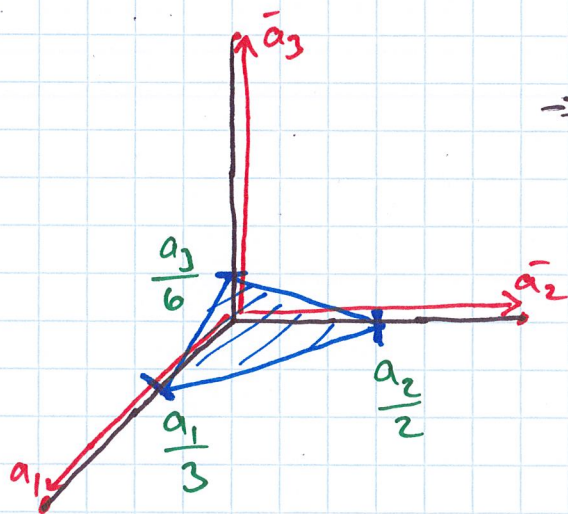
$$\therefore h:k:l = \frac{1}{x_1} : \frac{1}{x_2} : \frac{1}{x_3} = \frac{1}{2} : \frac{1}{3} : \frac{1}{1}$$

Find LCF for integer h, k, l :

$$h:k:l = 6 \times \left(\frac{1}{2} : \frac{1}{3} : \frac{1}{1} \right) = 3:2:6$$

\therefore Miller plane is: $(3, 2, 6)$

\rightarrow Typically Miller planes are drawn in the unit cell



\rightarrow equiv. plane family to original.